

## REGULARIZATION OF A DISCRETE BACKWARD PROBLEM USING COEFFICIENTS OF TRUNCATED LAGRANGE POLYNOMIALS

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ABSTRACT. We consider the problem of finding the initial temperature  $u(x, 0)$ , from a countable set of measured values  $\{u(x_j, 1)\}$ . The problem is severely ill-posed and a regularization is in order. Using the Hermite polynomials and coefficients of truncated Lagrange polynomials, we shall change the problem into an analytic interpolation problem and give explicitly a stable approximation. Error estimates and some numerical examples are given.

### 1. INTRODUCTION

Let  $u = u(x, t)$  represent a temperature distribution satisfying the heat equation

$$u_t - \Delta u = 0 \quad (x, t) \in \mathbb{R} \times (0, 1). \quad (1.1)$$

The backward problem is of finding the initial temperature  $u(x, 0)$  from the final temperature  $u(x, T)$ . For simplicity, we shall assume that  $T = 1$ . This is an ill-posed problem and has a long history [3]. This problem has been considered by many authors, using different approaches. The problem was studied intensively by the semi-group method associated with the quasi-reversibility method and the quasi-boundary value method; see for example [1, 2, 5, 7, 15, 22, 23, 12, 16, 13, 9, 26]. Using the Green function, we can transform the heat equation into

$$u(x, t) = \frac{1}{2\sqrt{\pi t}} \int_{-\infty}^{\infty} u(\xi, 0) e^{-\frac{(x-\xi)^2}{4t}} d\xi, \quad x \in \mathbb{R}, t > 0.$$

Hence

$$\frac{1}{\sqrt{\pi}} \int_{-\infty}^{\infty} u(2\xi, 0) e^{-(x-\xi)^2} d\xi = u(2x, 1).$$

In this form, we can consider the backward problem as the inversion Gaussian convolution (or Weierstrass transform) problem of finding  $u(2x, 0)$  from its image  $u(2x, 1)$ . Many inversion formulae for the Gauss transform were given in [18, 19, 20, 21]. In [13], using the reproducing kernel theory, the authors gave analytical inversion formulas which is optimal in an appropriate sense. In the latter paper, the case of nonexact  $L^2$ -data was studied and some sharp error estimates were given.

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Very recently, in [14], using the Paley-Wiener space and sinc approximation, the authors established a powerful practical numerical and analytical inversion formulas for the Gaussian convolution that is realized by computers. In [6, 27], the inversion Weierstrass transform for generalized functions was studied.

In practical situations, we get temperature measurements only at a discrete set of points, i.e.

$$u(x_j, 1) = \mu_j. \quad (1.2)$$

So, the problem of finding the initial temperature from discrete final values is necessary. In this case, the problem is severely ill-posed. Hence, a regularization is in order. However, the literature on this direction is very scarce. In [17], the authors used the shifted-Legendre polynomial to regularize a discrete form of the backward problem on the plane. However, the assumption that the temperature  $u(x, y)$  is of order  $e^{-(x^2+y^2)\alpha(x,y)}$  ( $\lim_{x,y \rightarrow -\infty} \alpha(x, y) = +\infty$ ) is very restrictive. In the present paper, the condition is removed completely.

As discussed, in the present paper, we shall consider a discrete form of the inversion problem for the Weierstrass transform

$$Wv(x_j) \equiv \frac{1}{\sqrt{\pi}} \int_{-\infty}^{\infty} v(\xi) e^{-(\frac{x_j}{2}-\xi)^2} d\xi = \mu_j. \quad (1.3)$$

where  $v(\xi) = u(2\xi, 0)$ . For the rest of this paper, we shall denote by  $Wv$  the sequence  $(Wv(x_j))$ .

Before going to the content of our paper, we shall give some definitions. In this paper, we denote

$$L^2_\rho(\mathbb{R}) = \{f : \mathbb{R} \rightarrow \mathbb{R} : f \text{ is Lebesgue measurable and } e^{-x^2/2} f \in L^2(\mathbb{R})\}.$$

The latter space is a Hilbert space with the norm

$$\|f\| = \left( \int_{-\infty}^{\infty} |f(x)|^2 e^{-x^2} dx \right)^{1/2}$$

and the inner product

$$\langle f, g \rangle = \int_{-\infty}^{\infty} f(x)g(x)e^{-x^2} dx, \quad \text{for } f, g \in L^2_\rho(\mathbb{R}).$$

We also denote

$$\ell^\infty = \{\mu = (\mu_j) : \mu_j \in \mathbb{R}, \sup_j |\mu_j| < \infty\}$$

with the norm  $\|\mu\|_\infty = \sup_j |\mu_j|$ .

For  $R > 0$ , we denote  $B_R = \{z \in \mathbb{C} : |z| < R\}$  and  $C_R = \{z \in \mathbb{C} : |z| = R\}$ . We also denote by  $H^1(B_R)$  the Hardy space of functions

$$\psi(z) = \sum_{n=0}^{\infty} \alpha_n z^n$$

analytic on the disc  $B_R$  with the norm

$$\|\psi\|_{H^1(B_R)}^2 = \sum_{n=0}^{\infty} |\alpha_n R^n|^2 < \infty.$$

Using the Parseval equality, we can rewrite the latter norm in another form

$$\|\psi\|_{H^1(B_R)}^2 = \frac{1}{2\pi} \int_0^{2\pi} |\psi(Re^{i\theta})|^2 d\theta.$$

If  $|\psi(Re^{i\theta})| \leq M$  for every  $\theta \in [0, 2\pi]$  then the latter equality gives

$$\|\psi\|_{H^1(B_R)}^2 \leq M.$$

Let  $v$  be an exact solution of (1.3), we recall that a sequence of linear operator  $T_n : \ell^\infty \rightarrow L_\rho^2(\mathbb{R})$  is a regularization sequence (or a regularizer) of Problem (1.3) if  $(T_n)$  satisfies two following conditions (see, [10])

(R1) For each  $n$ ,  $T_n$  is bounded,

(R2)  $\lim_{n \rightarrow \infty} \|T_n(Wv) - v\| = 0$ .

The number “ $n$ ” is called the regularization parameter. From (R1), (R2), we can obtain

(R3) For  $\epsilon > 0$ , there exists the functions  $n(\epsilon)$  and  $\delta(\epsilon)$  such that  $\lim_{\epsilon \rightarrow 0} n(\epsilon) = \infty$ ,  $\lim_{\epsilon \rightarrow 0} \delta(\epsilon) = 0$  and that

$$\|T_{n(\epsilon)}(\mu) - v\| \leq \delta(\epsilon)$$

for every  $\mu \in \ell^\infty$  such that  $\|\mu - Wv\|_\infty < \epsilon$ .

The number  $\epsilon$  is the error between the exact data  $Wv$  and the measured data  $\mu$ . For a given error  $\epsilon$ , there are infinitely many ways of choosing the regularization parameter  $n(\epsilon)$ . In the present paper, we give an explicit form of  $n(\epsilon)$ .

The remainder of the paper is divided into three sections. In Section 2, we shall transform the problem into an analytic interpolation problem and prove a uniqueness result. In Section 3, we shall find regularization functions by an association between Hermite polynomials and coefficients of Lagrange polynomials. Finally, in Section 4, some numerical examples are given.

## 2. REFORMULATION OF THE PROBLEM AND THE UNIQUENESS

Using Hermite polynomials (see [4, P. 65]) we can write

$$e^{-(z-\xi)^2} = \sum_{n=0}^{\infty} \frac{1}{n!} e^{-\xi^2} H_n(\xi) z^n,$$

where we recall that

$$H_n(\xi) = (-1)^n e^{\xi^2} \frac{d^n}{d\xi^n} e^{-\xi^2},$$

$$\langle H_n, H_m \rangle = \delta_{mn} \sqrt{\pi} 2^n n!$$

where  $\delta_{mn} = 0$  when  $n \neq m$  and  $\delta_{nn} = 1$ . We shall find a sequence  $(a_n)$  such that

$$v(\xi) = u(2\xi, 0) = \sum_{n=0}^{\infty} a_n H_n(\xi)$$

satisfies (1.3). From the orthogonality of  $\{H_n\}$  in the space  $L_\rho^2(R)$ , we can substitute the latter expansion into (1.3) to get

$$\mu_j = \sum_{n=0}^{\infty} a_n x_j^n.$$

Now, if we put

$$\phi(v)(z) = \sum_{n=0}^{\infty} a_n z^n, \tag{2.1}$$

then we have

$$\phi(v)(x_j) = \mu_j. \quad (2.2)$$

Hence, the problem is reformulated to the classical one of finding the sequence  $(a_n)$  (and of constructing a function  $v$ ) from the prescribed values  $(\mu_j)$  such that  $\phi(v)(z)$  satisfies (2.2). We first give some properties of the function  $\phi(v)$ .

**Lemma 2.1.** *Let  $v(x) = u(2x, 0)$  be in  $L^2_\rho(\mathbb{R})$ . If  $v$  has the expansion*

$$v(\xi) = \sum_{n=0}^{\infty} a_n H_n(\xi)$$

then

$$\sqrt{\pi} \sum_{n=0}^{\infty} |a_n|^2 2^n n! < \infty \quad (2.3)$$

and that the function  $\phi(v)(\cdot)$  is an entire function of order  $\rho \leq 2$ . Here we recall that the order of an entire function  $f$  is the number

$$\rho = \limsup_{r \rightarrow \infty} \frac{\ln \ln M_f(r)}{\ln r}$$

where  $M_f(r) = \max_{|z|=r} |f(z)|$ .

*Proof.* As mentioned before,  $\langle H_n, H_m \rangle = \delta_{mn} \sqrt{\pi} 2^n n!$  where  $\delta_{mn} = 0$  for  $m \neq n$  and  $\delta_{mm} = 1$ . Since

$$v(\xi) = \sum_{n=0}^{\infty} a_n H_n(\xi)$$

we get

$$\sqrt{\pi} \sum_{n=0}^{\infty} |a_n|^2 2^n n! = \|v\|^2 < \infty. \quad (2.4)$$

Now we prove that  $\phi(v)$  is an entire function. In fact, we consider the power series

$$\phi(v)(z) = \sum_{n=0}^{\infty} a_n z^n.$$

From (2.4), one has

$$|a_n|^2 \leq \frac{\|v\|^2}{2^n n!}.$$

It follows that

$$\lim_{n \rightarrow \infty} \sqrt[n]{|a_n|} = 0.$$

Hence, the power series has the convergent radius  $R = \infty$ , i.e.,  $\phi$  is an entire function. Now, we estimate the order of the entire function  $\phi$ . We note that the order  $\rho$  of  $\phi$  can be calculated by the following formula (see [11, P. 6])

$$\rho = \limsup_{n \rightarrow \infty} \frac{n \ln n}{\ln(1/|a_n|)}.$$

From (2.4), one has

$$1/|a_n|^2 \geq C 2^n n!$$

where  $C = \|v\|^{-2}$ . On the other hand, we have the Stirling formula (see [24, P. 688])

$$n! = \sqrt{2\pi n} n^n e^{-n} e^{\theta_n},$$

where

$$\frac{1}{12(n+1)} \leq \theta_n \leq \frac{1}{12(n-1)}.$$

Hence

$$2^n n! \geq \sqrt{2\pi n} (2/e)^n n^n.$$

It follows that

$$\ln(1/|a_n|^2) \geq C_1(1 + \ln \ln n + n \ln(2/e)) + n \ln n$$

where  $C_1$  is a generic constant. Hence

$$\rho = \limsup_{n \rightarrow \infty} \frac{2n \ln n}{\ln(1/|a_n|^2)} \leq \limsup_{n \rightarrow \infty} \frac{2n \ln n}{C_1(\ln \ln n + n \ln(2/e)) + n \ln n} = 2.$$

This completes the proof of Lemma 2.1.  $\square$

Now we have a uniqueness result.

**Theorem 2.2.** *Let  $\delta > 0$ . If*

$$\sum_{n=1}^{\infty} \frac{1}{|x_n|^{2+\delta}} = \infty$$

*then Problem (1.3) has at most one solution  $v \in L^2_\rho(\mathbb{R})$ .*

The latter condition means that the sequence  $(x_n)$  has an accumulation point on the extended real axis  $\mathbb{R} \cup \{\pm\infty\}$ . Moreover, if the accumulation point is  $\infty$  then the sequence  $(x_n)$  has to be “dense enough” near  $\infty$ .

*Proof.* Let  $v_1, v_2 \in L^2_\rho(\mathbb{R})$  be two solutions of (1.3). Putting  $v = v_1 - v_2$  and assuming that  $v = \sum_{n=1}^{\infty} a_n H_n$ , we shall get as in the beginning of Section 2

$$\phi(v)(x_j) = 0, \quad j = 1, 2, \dots$$

where  $\phi(v) = \sum_{n=1}^{\infty} a_n z^n$ . It follows that  $x_j$ 's are zeroes of the entire function  $\phi$ . If  $x_j$ 's has a finite accumulation point then the identity theorem shows that  $\phi(v) \equiv 0$ . If  $x_j$ 's do not have any finite accumulation points, we can assume, without loss of generality, that  $|x_1| \leq |x_2| < \dots$  and  $\lim_{j \rightarrow \infty} |x_j| = \infty$ . Since the order of  $\phi(v)$  is  $\leq 2$ , we get (see [11, P. 18])

$$\inf \left\{ \lambda \left| \sum_{n=1}^{\infty} \frac{1}{|x_n|^\lambda} < \infty \right. \right\} \leq \rho \leq 2.$$

It follows that

$$\sum_{n=1}^{\infty} \frac{1}{|x_n|^{2+\delta}} < \infty$$

which is a contradiction. Hence, in either cases, we have  $\phi(v) \equiv 0$ . It follows that  $a_n = 0$ ,  $n = 1, 2, \dots$ . This completes the proof.  $\square$

## 3. REGULARIZATION AND ERROR ESTIMATE

For the rest of this article, we shall assume that there exists an  $R > 0$  such that  $\sup_j |x_j| < R$ . Put  $\omega_n(z) = (z - x_0) \dots (z - x_n)$  and  $\mu = (\mu_j) \in \ell^\infty$ . We denote by  $L_n(\mu)$  the Lagrange polynomial of degree (at most)  $n$ , i.e.,

$$L_n(\mu)(z) = \sum_{j=0}^n \mu_j \frac{\omega_n(z)}{\omega_n'(x_j)(z - x_j)}$$

which satisfies  $L_n(\mu)(x_j) = \mu_j$ . Now, we denote by  $l_j^{(n)}(\mu)$  the coefficient of  $z^j$  in the expansion of the Lagrange polynomial  $L_n(\mu)$ , i.e.

$$L_n(z)(\mu) = \sum_{j=0}^n l_j^{(n)}(\mu) z^j. \quad (3.1)$$

We shall construct a regularization sequence. We denote by  $k_{0n}$  the greatest integer satisfying

$$n \ln \left( \frac{3}{2} \right) > (2k_{0n} + 1) \ln k_{0n}. \quad (3.2)$$

We can verify easily that  $\lim_{n \rightarrow \infty} k_{0n} = \infty$ . We choose a sequence  $(k_n)$  such that

$$0 < k_n \leq k_{0n}, \quad \lim_{n \rightarrow \infty} k_n = \infty. \quad (3.3)$$

For each  $n$ , we shall approximate the function  $v(x) = u(2x, 0)$  by the function

$$T_n(\mu)(x) = \sum_{j=0}^{k_n} l_j^{(n)}(\mu) H_j(x). \quad (3.4)$$

We shall verify that  $T_n$  is a regularization sequence. We first note that  $T_n : \ell^\infty \rightarrow L_\rho^2(\mathbf{R})$  is bounded, i.e., Condition (R1) (in Section 1) holds. In Theorem 3.1 below, we shall prove that  $(T_n)$  satisfies (R2) and, in Theorem 3.6, we shall prove that  $(T_n)$  satisfies (R3). In fact, we get the following regularization result for the case of exact data.

**Theorem 3.1.** *Let  $(k_n)$  be as in (3.3), let  $R \geq 1$  and let  $v \in L_\rho^2(\mathbb{R})$  be as in Theorem 2.2. Put  $F_n = T_n(Wv)$ . Then  $\|v - F_n\| \rightarrow 0$  as  $n \rightarrow \infty$ . Moreover, if  $v' \in L_\rho^2(\mathbb{R})$  then we can find an  $n_0$  such that*

$$\|v - F_n\|^2 \leq \|v\|^2 e^{8R^2} \left( \frac{2}{3} \right)^n + \frac{1}{k_n} \|v'\|^2 \quad \text{for } n \geq n_0.$$

*If we choose  $k_n = k_{0n}$ ,  $n = 1, 2, \dots$  then the latter inequality can be rewritten as follows*

$$\|v - F_n\|^2 \leq \|v\|^2 e^{8R^2} \left( \frac{2}{3} \right)^n + \frac{1}{\sqrt{n}} \|v'\|^2 \quad \text{for } n \geq n_0.$$

Before proving this theorem, some remarks are in order. We note that the coefficients of  $z^j$  ( $j \geq k_n + 1$ ) in the expansion of the Lagrange polynomial (3.1) are truncated in (3.4). If we use coefficients of  $z^j$ 's (for  $j$  large) of  $L_n$  in (3.4) then we shall get functions which are unstable approximation of  $v$ . To illustrate the latter fact, in Section 4, we shall give a numerical example. In fact, we can say that the polynomial

$$L_{nk_n}(z) = \sum_{j=0}^{k_n} l_j^{(n)} z^j$$

is a truncated Lagrange polynomial (see [25] for a similar definition). Hence, our method of regularization is of using the coefficients of truncated Lagrange polynomials. We shall give an estimate for  $l_j^{(n)}$  and the proof of Theorem 3.1. To this end, some lemmas will be established.

**Lemma 3.2.** *Let  $v, \phi(v)$  be as in Lemma 2.1. Then  $\phi : L^2_\rho(\mathbb{R}) \rightarrow H^1(B_R)$  is a bounded linear operator satisfying*

$$\|\phi(v)\|_{H^1(B_R)}^2 \leq e^{R^2/2} \|v\|^2.$$

*Proof.* We have

$$\begin{aligned} \|\phi(v)\|_{H^1(B_R)}^2 &= \sum_{n=0}^{\infty} |a_n|^2 R^{2n} \leq \sqrt{\pi} \sum_{n=0}^{\infty} |a_n|^2 n 2^n n! \frac{R^{2n}}{n 2^n n!} \\ &\leq \|v\|^2 \sum_{n=0}^{\infty} \frac{R^{2n}}{2^n n!} \\ &= e^{R^2/2} \|v\|^2. \end{aligned}$$

This completes the proof  $\square$

**Lemma 3.3.** *Let  $v, \phi(v)$  and  $(a_n)$  be as in Lemma 2.1. Assume that  $(x_j)$  is in the disc  $B_R$ . Then one has*

$$\sum_{j=0}^n R^{2j} |a_j - l_j^{(n)}|^2 + \sum_{j=n+1}^{\infty} R^{2j} |a_j|^2 \leq \frac{1}{9} \left(\frac{2}{3}\right)^{2n} e^{8R^2} \|v\|^2.$$

*Proof.* In the present proof, we shall denote  $\phi(v)$  by  $\phi$ . We have the Hermitian representation (see [8, P. 58])

$$\phi(z) - L_n(z) = \frac{1}{2\pi i} \int_{C_{4R}} \frac{\omega_n(z)}{\omega_n(t)} \cdot \frac{\phi(t)}{t-z} dt.$$

Now, for every  $t \in C_{4R}$  one has  $|\omega_n(t)| \geq (3R)^n$ . On the other hand, one has for every  $|z| \leq R$

$$|\omega_n(z)| \leq (2R)^n.$$

We claim that

$$\|\phi - L_n\|_{H^1(B_R)} \leq \frac{1}{3} \left(\frac{2}{3}\right)^n \|\phi\|_{H^1(B_{4R})}.$$

In fact, we have for  $|z| = R, t = 4Re^{i\theta}$

$$\begin{aligned} |\phi(z) - L_n(z)| &\leq \frac{1}{2\pi} \int_0^{2\pi} \frac{(2R)^n}{(3R)^n} \cdot \frac{|\phi(4Re^{i\theta})|}{4R-R} R d\theta \\ &\leq \frac{1}{3} \left(\frac{2}{3}\right)^n \frac{1}{2\pi} \int_0^{2\pi} |\phi(4Re^{i\theta})| d\theta \\ &\leq \frac{1}{3} \left(\frac{2}{3}\right)^n \left(\frac{1}{2\pi} \int_0^{2\pi} |\phi(4Re^{i\theta})|^2 d\theta\right)^{1/2} \\ &= \frac{1}{3} \left(\frac{2}{3}\right)^n \|\phi\|_{H^1(B_{4R})}. \end{aligned}$$

It follows that

$$\|\phi - L_n\|_{H^1(B_R)}^2 = \frac{1}{2\pi} \int_0^{2\pi} |\phi(Re^{i\theta}) - L_n(Re^{i\theta})|^2 d\theta \leq \frac{1}{9} \left(\frac{2}{3}\right)^{2n} \|\phi\|_{H^1(B_{4R})}^2$$

as claimed. In view of Lemma 3.2, it follows that

$$\sum_{j=0}^n R^{2j} |a_j - l_j^{(n)}|^2 + \sum_{j=n+1}^{\infty} R^{2j} |a_j|^2 \leq \frac{1}{9} \left(\frac{2}{3}\right)^{2n} e^{8R^2} \|v\|^2.$$

This completes the proof.  $\square$

**Lemma 3.4.** *Let  $f \in L^2_\rho(\mathbb{R})$  satisfy  $f' \in L^2_\rho(\mathbb{R})$  and  $f = \sum_{n=0}^{\infty} c_n H_n$ . Then we have*

$$\sum_{n=0}^{\infty} 2nc_n^2 \sqrt{\pi} 2^n n! = \|f'\|^2.$$

*Proof.* We note that  $H_n$  satisfies the differential equation

$$y'' - 2xy' + 2ny = 0,$$

(see [4, P. 66]). It follows that  $H_n$  satisfies

$$(e^{-x^2} y')' + 2nye^{-x^2} = 0.$$

Hence we have

$$(e^{-x^2} f')' = \sum_{n=0}^{\infty} c_n (e^{-x^2} H_n')' = \sum_{n=0}^{\infty} -2nc_n H_n e^{-x^2}.$$

Hence, taking the inner product in  $L^2(\mathbb{R})$  with respect  $H_n$ , we get in view of the orthogonality

$$\int_{-\infty}^{\infty} e^{-x^2} f'(x) c_n H_n'(x) dx = 2nc_n^2 \sqrt{\pi} 2^n n!.$$

It follows that

$$\int_{-\infty}^{\infty} e^{-x^2} f'(x) f'(x) dx = \sum_{n=0}^{\infty} 2nc_n^2 \sqrt{\pi} 2^n n!.$$

This completes the proof.  $\square$

**Lemma 3.5.** *For  $(k_{0n})$  as in (3.2), there exist  $a_0, n_0 > 0$  such that*

$$\left(\frac{3}{2}\right)^n \geq \sqrt{\pi} j! 2^j \quad \text{for } 0 \leq j \leq k_{0n}$$

*for every  $n > a_0$  and that  $k_{0n} \geq \sqrt{n}$  for every  $n > n_0$ .*

*Proof.* For every  $k > 4\pi^2 e^2$ , we have

$$\begin{aligned} \ln \left( 2\pi e \sqrt{k} (2ek)^k \right) &= \ln(2\pi e) + \frac{1}{2} \ln k + k(1 + \ln 2 + \ln k) \\ &\leq \ln(2\pi e) + k(1 + \ln 2) + \frac{1}{2} \ln k + k \ln k \\ &\leq (2k + 1) \ln k \equiv g(k) \end{aligned}$$

For every  $n > a_0 = g(576) \ln^{-1} \left(\frac{3}{2}\right)$ , one has in view of the definition of  $k_{0n}$  that  $k_{0n} \geq 576 > 4\pi^2 e^2$ . Hence, we have for  $n > a_0$

$$(2k_{0n} + 1) \ln k_{0n} \geq \ln \left( 2\pi e \sqrt{k_{0n}} (2ek_{0n})^{k_{0n}} \right).$$

Now, since  $k_{0n}$  satisfies

$$n \ln \left(\frac{3}{2}\right) > (2k_{0n} + 1) \ln k_{0n},$$

we have for  $n > a_0$

$$n \ln \left( \frac{3}{2} \right) > \ln \left( 2\pi e \sqrt{k_{0n}} (2ek_{0n})^{k_{0n}} \right).$$

Using Stirling formula we get

$$\sqrt{\pi} k_{0n}! 2^{k_{0n}} \leq 2\pi e \sqrt{k_{0n}} (2ek_{0n})^{k_{0n}}.$$

It follows that

$$\left( \frac{3}{2} \right)^n > 2\pi e \sqrt{k_{0n}} (2ek_{0n})^{k_{0n}} \geq \sqrt{\pi} k_{0n}! 2^{k_{0n}}.$$

Since

$$(2k_{0n} + 3) \ln(k_{0n} + 1) \geq n \ln \left( \frac{3}{2} \right) > (2k_{0n} + 1) \ln k_{0n}$$

one has  $k_{0n} \rightarrow \infty$  as  $n \rightarrow \infty$  and that

$$\lim_{n \rightarrow \infty} \frac{n \ln \left( \frac{3}{2} \right)}{k_{0n} \ln k_{0n}} = 2.$$

It follows that

$$\lim_{n \rightarrow \infty} \frac{\sqrt{n}}{k_{0n}} = \lim_{n \rightarrow \infty} \frac{\sqrt{n}}{\sqrt{2k_{0n} \ln k_{0n}}} \cdot \frac{\sqrt{2k_{0n} \ln k_{0n}}}{k_{0n}} = 0.$$

Hence, we can find an  $n_0 > a_0$  such that  $k_{0n} \geq \sqrt{n}$  for every  $n \geq n_0$ . This completes the proof.  $\square$

*Proof of Theorem 3.1.* For  $R \geq 1$ , it follows in view of the orthogonality of  $(H_n)$  that

$$\begin{aligned} \|v - F_n\|^2 &= \sum_{j=0}^{k_n} |a_j - l_j^{(n)}|^2 \sqrt{\pi} j! 2^j + \sum_{j=k_n+1}^{\infty} |a_n|^2 \sqrt{\pi} j! 2^j \\ &= \sum_{j=0}^{k_n} R^{2j} |a_j - l_j^{(n)}|^2 \frac{\sqrt{\pi} j! 2^j}{R^{2j}} + \sum_{j=k_n+1}^{\infty} |a_j|^2 \sqrt{\pi} j! 2^j \\ &\leq \sum_{j=0}^{k_n} R^{2j} |a_j - l_j^{(n)}|^2 \sqrt{\pi} j! 2^j + \sum_{j=k_n+1}^{\infty} |a_j|^2 \sqrt{\pi} j! 2^j. \end{aligned}$$

Using Lemma 3.3, we have

$$\|v - F_n\|^2 \leq \|v\|^2 e^{8R^2} \frac{1}{9} \left( \frac{2}{3} \right)^{2n} \sqrt{\pi} k_n! 2^{k_n} + \sum_{j=k_n+1}^{\infty} |a_j|^2 \sqrt{\pi} j! 2^j.$$

In view of Lemma 3.5,

$$\|v - F_n\|^2 \leq \|v\|^2 e^{8R^2} \left( \frac{2}{3} \right)^n + \sum_{j=k_n+1}^{\infty} |a_j|^2 \sqrt{\pi} j! 2^j.$$

Using Lemma 2.1, one gets

$$\lim_{n \rightarrow \infty} \sum_{j=k_n+1}^{\infty} |a_j|^2 \sqrt{\pi} j! 2^j = 0.$$

It follows that

$$\lim_{n \rightarrow \infty} \|v - F_n\| = 0.$$

Now, if  $v' \in L^2_\rho(\mathbb{R})$  then we get in view of Lemma 3.4

$$\sum_{j=k_n+1}^{\infty} |a_j|^2 \sqrt{\pi} j! 2^j \leq \frac{1}{k_n} \|v'\|^2,$$

which gives

$$\|v - F_n\|^2 \leq \|v\|^2 \left(\frac{2}{3}\right)^n + \frac{1}{k_n} \|v'\|^2.$$

This proves the first estimate of Theorem 3.1. Now if  $k_n = k_{0n}$  then Lemma 3.5 gives  $k_n \geq \sqrt{n}$  for  $n \geq n_0$ . Hence

$$\|v - F_n\|^2 \leq \|v\|^2 \left(\frac{2}{3}\right)^n + \frac{1}{\sqrt{n}} \|v'\|^2.$$

This completes the proof.  $\square$

Now, we consider the case of nonexact data. Let  $\epsilon > 0$  and let  $\mu^\epsilon = (\mu_j^\epsilon)$  be a nonexact data of  $(Wv(x_j)) = (u(x_j, 1))$  satisfying

$$\sup_j |u(x_j, 1) - \mu_j^\epsilon| < \epsilon.$$

We first put

$$D_m = \max_{1 \leq n \leq m} \left( \max_{|z| \leq R} \left| \frac{\omega_m(z)}{(z - x_n)\omega'_m(x_n)} \right| \right)$$

and

$$F_n^\epsilon = T_n(\mu^\epsilon) = \sum_{j=0}^{k_n} l_{j^\epsilon}^{(n)} H_j,$$

where  $l_{j^\epsilon}^{(n)}$  is the coefficient of  $z^j$  in the expansion of the Lagrange polynomial

$$L_{en} = \sum_{j=0}^n \mu_j^\epsilon \frac{\omega_n(z)}{\omega'_n(z_j)(z - z_j)}.$$

Let  $\psi$  be an increasing function such that

$$\psi(n) \geq (n+1)D_n \left(\frac{3}{2}\right)^{n/2}, \quad \lim_{x \rightarrow \infty} \psi(x) = \infty$$

and

$$n(\epsilon) = [\psi^{-1}(\epsilon^{-\frac{1}{2}})] + 1$$

where  $[x]$  is the greatest integer  $\leq x$ . Using the latter function, we shall prove that  $(T_n)$  satisfies the condition (R3).

**Theorem 3.6.** *Let  $R > 1$  and let  $v \in L^2_\rho(\mathbb{R})$ . Let  $\epsilon > 0$  and let  $(\mu_j^\epsilon)$  be a measured data of  $(u(x_j, 1))$  satisfying*

$$\sup_j |u(x_j, 1) - \mu_j^\epsilon| < \epsilon.$$

*Then*

$$\|v - F_{n(\epsilon)}^\epsilon\| \leq \delta(\epsilon) = \|v - F_{n(\epsilon)}\| + \sqrt{\epsilon}.$$

*Moreover, if  $v' \in L^2_\rho(\mathbb{R})$  then*

$$\|v - F_{n(\epsilon)}^\epsilon\|^2 \leq 2\|v\|^2 \left(\frac{2}{3}\right)^{n(\epsilon)} + \frac{2}{k_{n(\epsilon)}} \|v'\|^2 + 2\epsilon.$$

In the latter inequality, if  $k_n = k_{0n}$  then there exists an  $\epsilon_0 > 0$  such that

$$\|v - F_{n(\epsilon)}^\epsilon\|^2 \leq 2\|v\|^2 \left(\frac{2}{3}\right)^{n(\epsilon)} + \frac{2}{\sqrt{n(\epsilon)}} \|v'\|^2 + 2\epsilon.$$

for  $0 < \epsilon < \epsilon_0$ .

*Proof.* We first claim that

$$\|F_n - F_n^\epsilon\| \leq (n+1) \left(\frac{3}{2}\right)^{n/2} \epsilon D_n.$$

In view of Lemma 3.5,

$$\begin{aligned} \|F_n - F_n^\epsilon\|^2 &= \sum_{j=0}^{k_n} |l_j^{(n)} - l_{j\epsilon}^{(n)}|^2 \sqrt{\pi} j! 2^j \\ &= \sum_{j=0}^{k_n} R^{2j} |l_j^{(n)} - l_{j\epsilon}^{(n)}|^2 \frac{\sqrt{\pi} j! 2^j}{R^{2j}} \\ &\leq \left(\frac{3}{2}\right)^n \sum_{j=0}^{k_n} R^{2j} |l_j^{(n)} - l_{j\epsilon}^{(n)}|^2. \end{aligned}$$

Hence

$$\|F_n - F_n^\epsilon\|^2 \leq \left(\frac{3}{2}\right)^n \|L_n - L_{n\epsilon}\|_{H^1(B_R)}^2. \quad (3.5)$$

On the other hand

$$L_n(z) - L_{n\epsilon}(z) = \sum_{j=0}^n (\mu_j - \mu_j^\epsilon) \frac{\omega_n(z)}{\omega_n'(x_j)(z - x_j)}.$$

It follows that

$$\|L_n - L_{n\epsilon}\|_{H^1(B_{1R})} \leq \sum_{j=0}^n |\mu_j - \mu_j^\epsilon| D_n \leq (n+1) \epsilon D_n.$$

So that

$$\|F_n - F_n^\epsilon\| \leq (n+1) \left(\frac{3}{2}\right)^{n/2} \epsilon D_n.$$

Now, we have

$$\|v - F_n^\epsilon\| \leq \|v - F_n\| + \|F_n^\epsilon - F_n\|.$$

Hence

$$\|v - F_n^\epsilon\| \leq \|v - F_n\| + \epsilon(n+1) D_n \left(\frac{3}{2}\right)^{n/2}.$$

For  $n = n(\epsilon)$ , we get in view of the definition of  $n(\epsilon)$  that

$$\|v - F_{n(\epsilon)}^\epsilon\| \leq \|v - F_{n(\epsilon)}\| + \sqrt{\epsilon}.$$

Since  $n(\epsilon) \rightarrow \infty$  as  $\epsilon \rightarrow 0$ , we can get from Theorem 3.1 and the latter inequality that

$$\lim_{\epsilon \rightarrow 0} \|v - F_{n(\epsilon)}^\epsilon\| = 0.$$

Now if  $v' \in L^2_\rho(\mathbb{R})$ , Theorem 3.1 and (3.5) give

$$\|v - F_n^\epsilon\|^2 \leq 2\|v\|^2 \left(\frac{2}{3}\right)^n + \frac{2}{k_n} \|v'\|^2 + 2(n+1)^2 \epsilon^2 D_n^2 \left(\frac{3}{2}\right)^n.$$

From the definition of  $n(\epsilon)$ , one has

$$\|v - F_{n(\epsilon)}^\epsilon\|^2 \leq 2\|v\|^2 \left(\frac{2}{3}\right)^{n(\epsilon)} + \frac{2}{k_{n(\epsilon)}} \|v'\|^2 + 2\epsilon.$$

Finally, if  $k_n = k_{0n}$  then Lemma 3.5 shows that, there exists an  $\epsilon_0 > 0$  such that  $k_{n(\epsilon)} \geq \sqrt{n(\epsilon)}$  for every  $0 < \epsilon < \epsilon_0$ . Hence, we shall get the desired estimate. This completes the proof.  $\square$

#### 4. NUMERICAL EXAMPLES

We shall give two numerical examples. In the first example, we consider  $x_j = \frac{1}{1+j}$ ,  $j = 0, 1, \dots, 100$ . We choose the exact function  $v(\xi) = 1$  and the non-exact data  $\mu_j^\epsilon = 1 + \frac{1}{2 \cdot 10^{20(j+1)}}$ . From the latter data, we can calculate (using MAPLE) the first six coefficients of the corresponding Lagrange polynomial  $[l_0^{(100)}, l_1^{(100)}, l_2^{(100)}, l_3^{(100)}, l_4^{(100)}, l_5^{(100)}]$  which are

$$s := [1 + 2575.000000 \times 10^{-20}, -6.546062500 \times 10^{-14}, 1.094478041 \times 10^{-10}, \\ -1.354054633 \times 10^{-7}, 1.322015356 \times 10^{-4}, -1.060903238 \times 10^{-1}].$$

Using the first five coefficients of the corresponding Lagrange polynomial, we get the approximation  $F_1 = \sum_{j=0}^4 l_j^{(100)} H_j$  of  $v$

$$F_1 := 1.001586418 + 0.000001624865429x - 0.006345673271x^2 \\ - 0.000001083243706x^3 + 0.002115224570x^4.$$

We have

$$\int_{-20}^{20} |F_1(x) - v(x)| e^{-x^2} dx \simeq 0.003448971524.$$

We have the graphs of two functions  $v$  and  $F_1$ . The approximation is very good in the interval  $[-2, 2]$ .

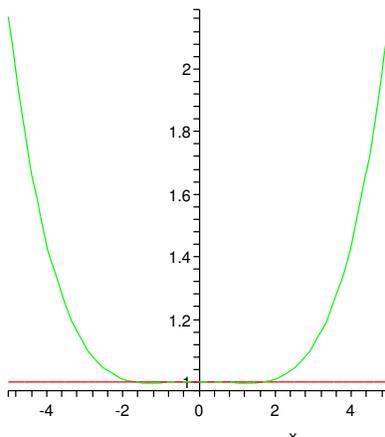


FIGURE 1. the graphs of  $v$  and  $F_1$  on  $[-4, 4]$

If we use the first six coefficients of the Lagrange polynomial, we get the approximation  $F_2 = \sum_{j=0}^5 l_j^{(100)} H_j$  of  $v$

$$F_2 := 1.001586418 - 12.73083724x - 0.006345673271x^2 \\ + 16.97445073x^3 + 0.002115224570x^4 - 3.394890362x^5.$$

We have

$$\int_{-20}^{20} |F_2(x) - v(x)| e^{-x^2} dx \simeq 8.752434897.$$

In this case, we can see that the error is larger than the foregoing case.

In the second example, we consider  $x_j = \frac{1}{1+j}$ ,  $j = 0, 1, \dots, 140$ . We choose the exact function  $v(\xi) = 1$ ,  $\mu_j^s = 1 + \frac{1}{2 \cdot 10^{20(j+1)}}$ . From the latter data, we can calculate the first six coefficients of Lagrange polynomial  $[l_0^{(140)}, l_1^{(140)}, l_2^{(140)}, l_3^{(140)}, l_4^{(140)}, l_5^{(140)}]$  which are

$$s := [1 + 5005.000000 \times 10^{-20}, -2.481893750 \times 10^{-13}, 8.126181478 \times 10^{-10}, \\ -0.1976424306 \times 10^{-5}, 0.3808576622 \times 10^{-2}, -6.056645660].$$

Using the first five coefficients of the corresponding Lagrange polynomial, we get the approximation  $F_3 = \sum_{j=0}^4 l_j^{(140)} H_j$  of  $v$

$$F_3 := 1.045702918 + 0.00002371709117x - 0.1828116746x^2 \\ - 0.00001581139445x^3 + 0.06093722595x^4.$$

We have

$$\int_{-20}^{20} |F_3(x) - v(x)| e^{-x^2} dx \simeq 0.09936096138.$$

On the other hand, if we use the first six coefficients of the Lagrange polynomial, we have the function  $F_4 = \sum_{j=0}^5 l_j^{(140)} H_j$

$$F_4 := 1.045702918 - 726.7974555x - 0.1828116746x^2 + 969.0632898x^3 \\ + 0.06093722595x^4 - 193.8126611x^5.$$

We have an error estimate

$$\int_{-20}^{20} |F_4(x) - v(x)| e^{-x^2} dx \simeq 499.6722779.$$

This case shows that the error is very large if we use too many coefficients of the Lagrange polynomial.

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